

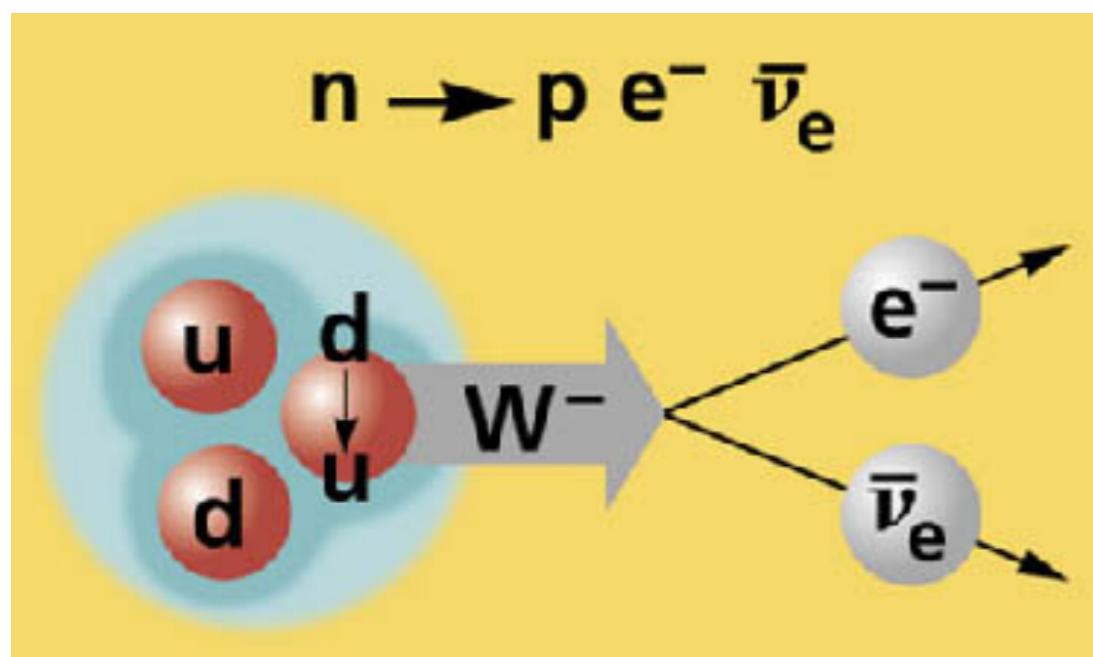
# Axial properties of N and N<sup>\*</sup> resonances

Ki-Seok Choi, W. Plessas, and R.F. Wagenbrunn

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Karl-Franzens-Universität Graz, Austria

# Weak Transitions



Decay parameter of n

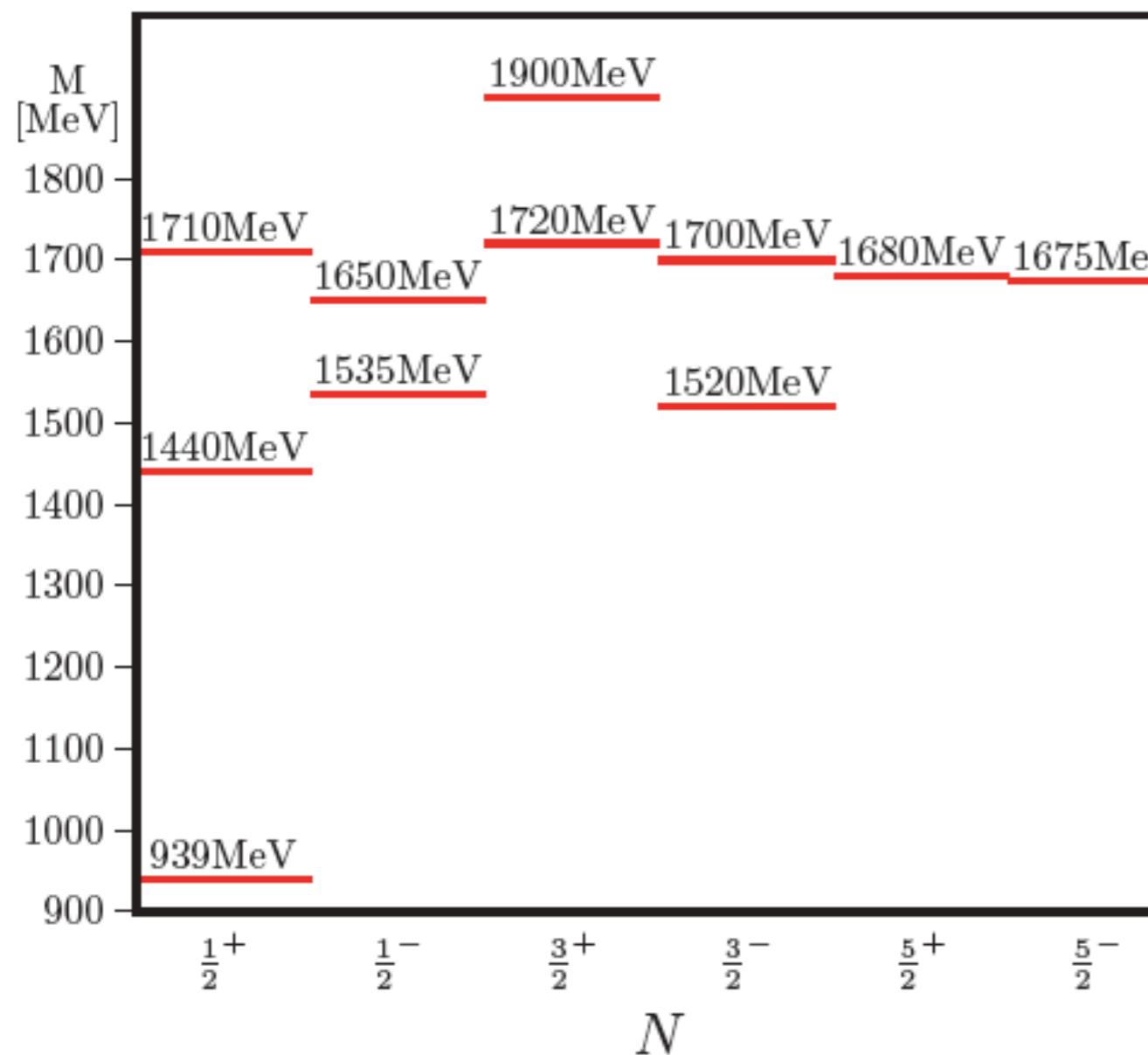
$$\lambda \equiv g_A/g_V = -1.2694 \pm 0.0028$$

Particle Data Group (PDG) 2010

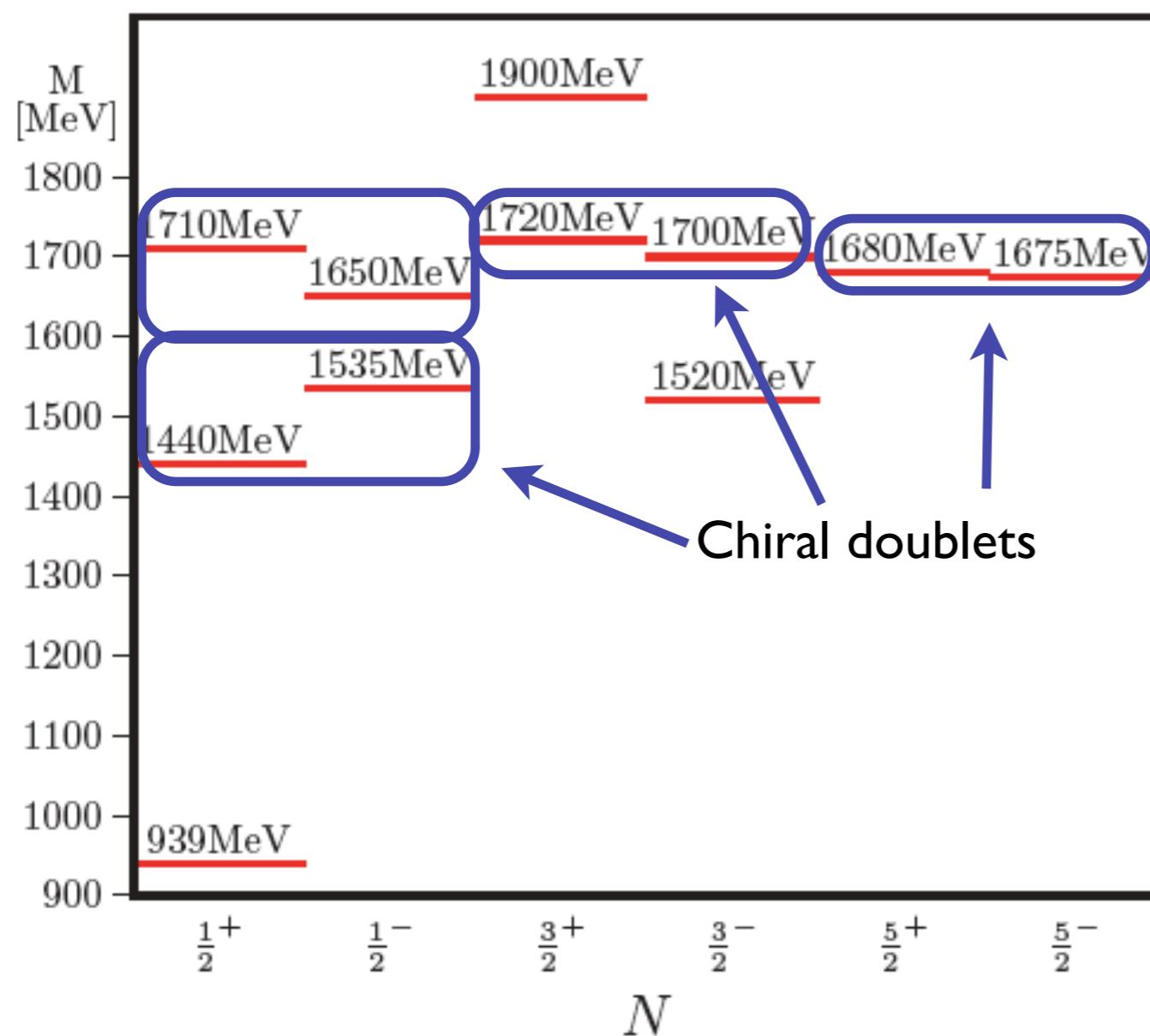
How about resonant states  $N^*$ ?

Recently, attracted much interest in studies  
of the strong interactions (QCD).

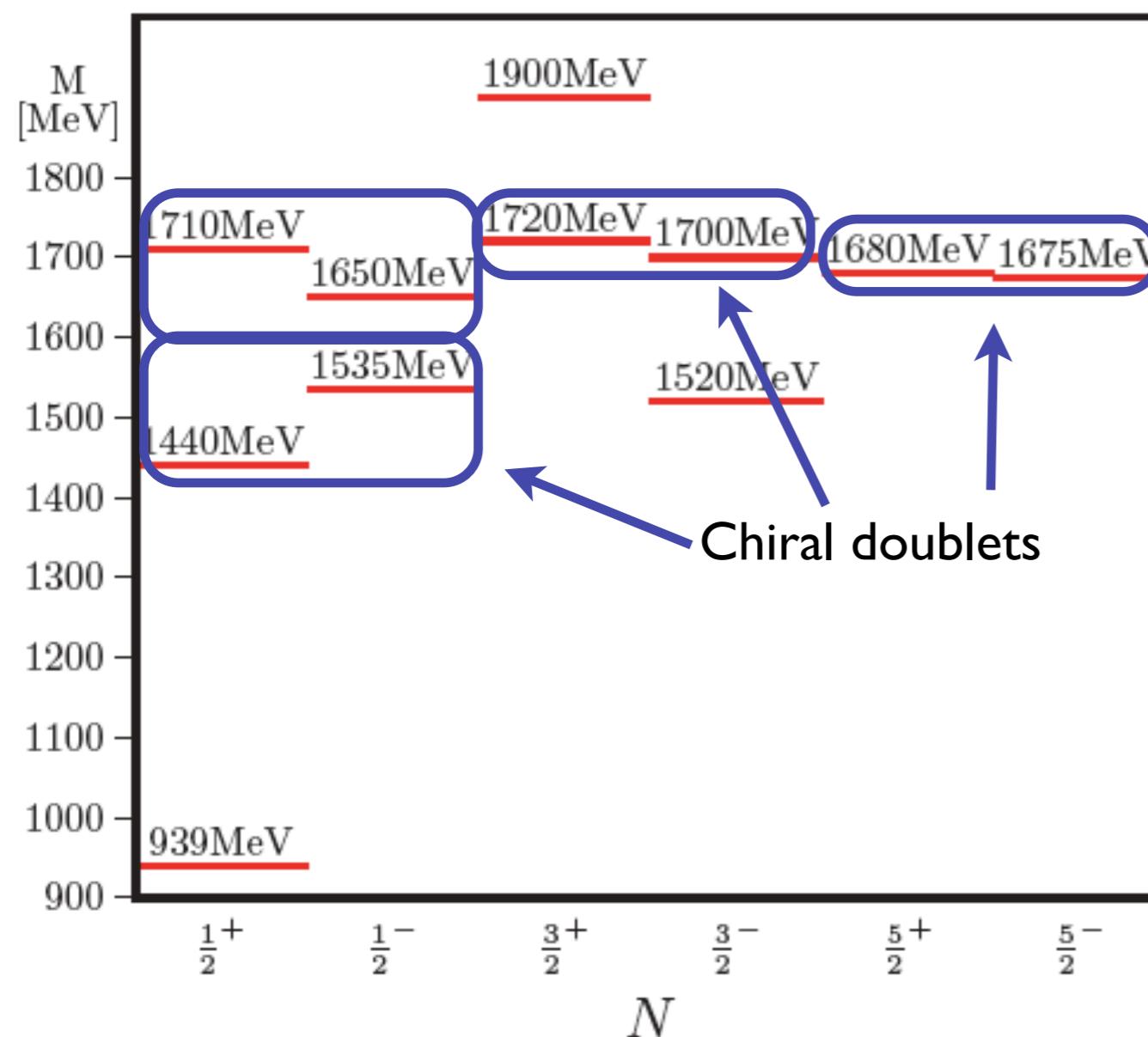
# Chiral Restoration Scenario of QCD



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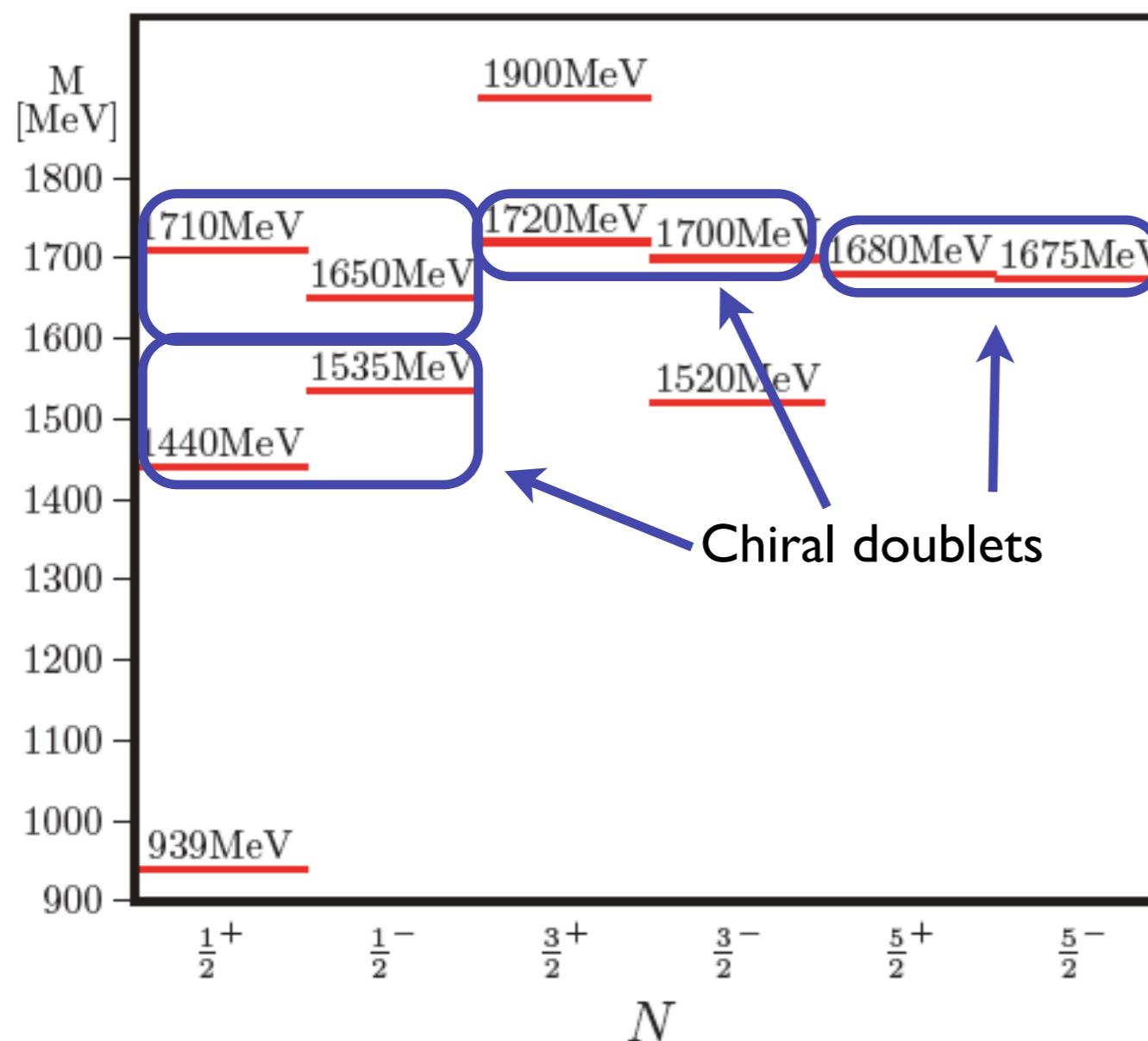


## Goldberger-Treiman Relation

$$g_{\pi NN} = \frac{g_A M_N}{f_\pi}$$

↑ Axial charge      ↑ Nucleon mass  
                         ↓ Pion Nucleon coupling      ↓ Pion decay constant

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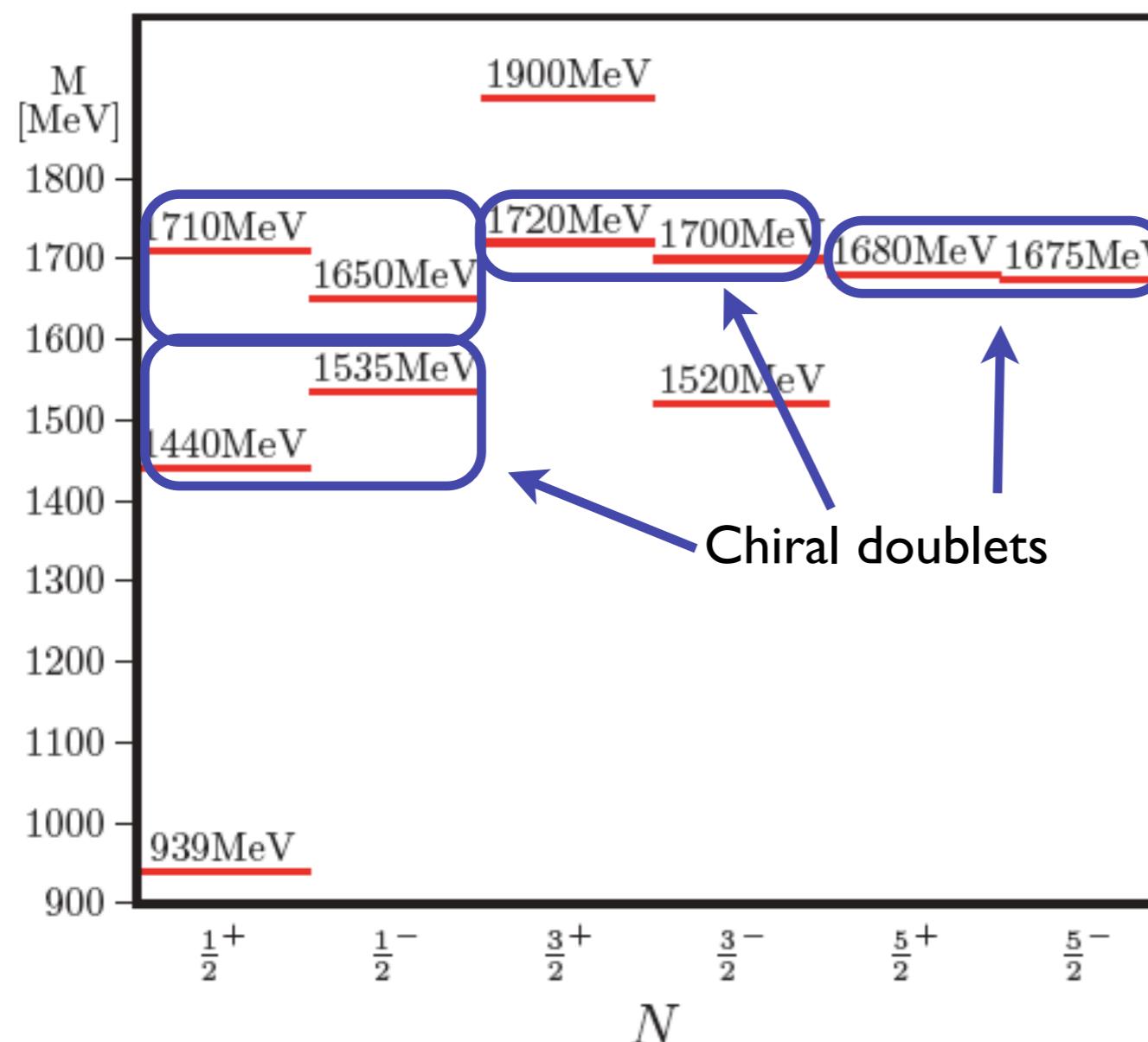
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Restoration of the chiral symmetry :

$$g_{\pi NN} \simeq 0$$

L.Ya. Glozman: Phys. Rev. Lett. 99, 191602 (2007)

# Chiral Restoration Scenario of QCD



## Goldberger-Treiman Relation

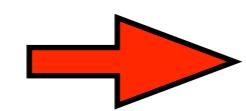
$$g_{\pi NN} = \frac{g_A M_N}{f_\pi}$$

Axial charge       Nucleon mass   
 Pion Nucleon coupling       Pion decay constant 

Restoration of the chiral symmetry :

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 Requirement for restoration :  $g_A \simeq 0$

# Analytic and Lattice Results

## Non-relativistic SU(6) wave function

$$\frac{1}{2}^+, N(939) : G^A = \frac{5}{3} \sim 1.66$$

$$\frac{1}{2}^+, N(1440) : G^A = \frac{5}{3} \sim 1.66$$

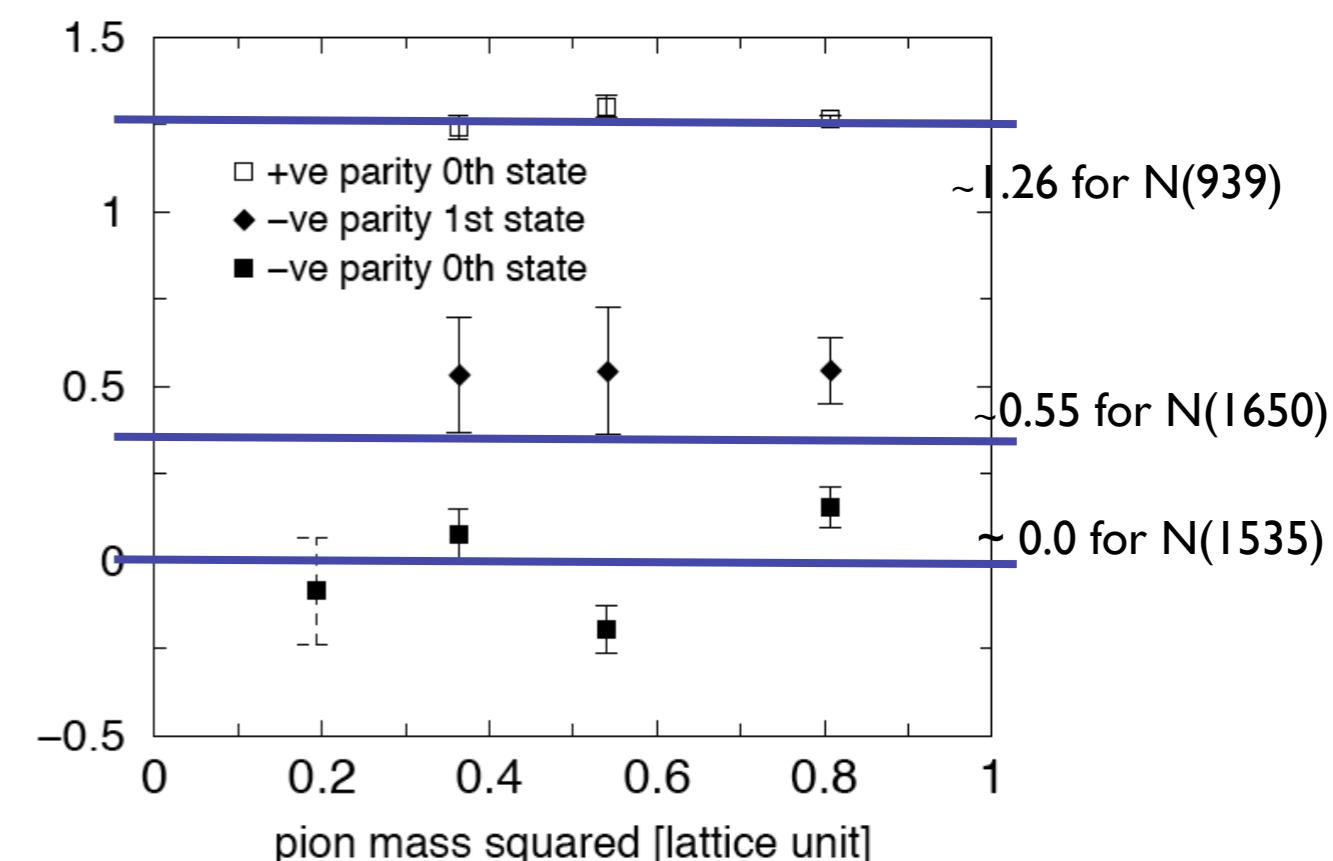
$$\frac{1}{2}^+, N(1710) : G^A = \frac{1}{3} \sim 0.33$$

$$\frac{1}{2}^-, N(1535) : G^A = -\frac{1}{9} \sim -0.11$$

$$\frac{1}{2}^-, N(1650) : G^A = \frac{5}{9} \sim 0.55$$

Small

## Lattice results



# Relativistic Constituent Quark Model

- Framework: Relativistic Quantum Mechanics (RQM)

i.e. quantum theory respecting Poincaré invariance  
 (theory on  $\mathcal{H}$  corresponding to a finite number of particles, not a field theory)

- Invariant mass operator

$$\hat{M} = \hat{M}_{free} + \hat{M}_{int} = \sqrt{\hat{H}_0^2 - \hat{\vec{P}}_{free}^2} + \sum_{i < j} [\hat{V}_{ij}^{conf} + \hat{V}_{ij}^{hf}]$$

Eigenvalue equation :  $\hat{M}|P, J, \Sigma\rangle = m|P, J, \Sigma\rangle$

Relativistic kinetic energy :  $H_0 = \sum_{i=1}^3 \sqrt{\vec{p}_i^2 + m_i^2}$

Confinement :  $\hat{V}_{ij}^{conf} = V_0 + C r_{ij}$

Hyperfine interaction  $\hat{V}_{ij}^{hf}$  : OGE, psGBE, EGBE

# Axial Form Factors

- Matrix elements of the axial current operator

$$\left\langle P, \frac{1}{2}, \Sigma' \left| \hat{A}_a^\mu \right| P, \frac{1}{2}, \Sigma \right\rangle = \bar{U}(P, \Sigma') g_A \gamma^\mu \gamma_5 \frac{\tau_a}{2} U(P, \Sigma)$$

$$\left\langle P, \frac{3}{2}, \Sigma' \left| \hat{A}_a^\mu \right| P, \frac{3}{2}, \Sigma \right\rangle = \bar{U}^\nu(P, \Sigma') g_A \gamma^\mu \gamma_5 \frac{\tau_a}{2} U_\nu(P, \Sigma)$$

$$\left\langle P, \frac{5}{2}, \Sigma' \left| \hat{A}_a^\mu \right| P, \frac{5}{2}, \Sigma \right\rangle = \bar{U}^{\nu\lambda}(P, \Sigma') g_A \gamma^\mu \gamma_5 \frac{\tau_a}{2} U_{\nu\lambda}(P, \Sigma)$$

$g_A$  : Axial charge of N and  $N^*$



Incoming / Outgoing baryon states :  $|V, M, J, \Sigma\rangle = |P, J, \Sigma\rangle$

$$\langle V, M, J, \Sigma | \hat{A}_a^\mu | V, M, J, \Sigma \rangle = \frac{2}{MM'} \sum_{\sigma_i \sigma'_i} \sum_{\mu_i \mu'_i} \int d^3 \vec{k}_2 d^3 \vec{k}_3 d^3 \vec{k}'_2 d^3 \vec{k}'_3$$

$$\times \sqrt{\frac{(\sum_i \omega'_i)^3}{\prod_i 2\omega'_i}} \prod_{\sigma'_i} D_{\sigma'_i \mu'_i}^{*\frac{1}{2}} \{ R_W[k'_i; B(V')] \} \Psi_{M'J'\Sigma'}^* \left( \vec{k}'_1, \vec{k}'_2, \vec{k}'_3; \mu'_1, \mu'_2, \mu'_3 \right) \langle p'_1, p'_2, p'_3; \sigma'_1, \sigma'_2, \sigma'_3 | \hat{A}_{a,rd}^\mu | p_1, p_2, p_3; \sigma_1, \sigma_2, \sigma_3 \rangle$$

$$\times \sqrt{\frac{(\sum_i \omega_i)^3}{\prod_i 2\omega_i}} \prod_{\sigma_i} D_{\sigma_i \mu_i}^{\frac{1}{2}} \{ R_W[k_i; B(V)] \} \Psi_{MJ\Sigma} \left( \vec{k}_1, \vec{k}_2, \vec{k}_3; \mu_1, \mu_2, \mu_3 \right) 2MV_0 \delta^3 \left( M\vec{V} - M'\vec{V}' - \vec{q} \right)$$

where  $p_i = B_c(V)k_i$ ,  $p'_i = B_c(V')k'_i$  and  $w_i = \sqrt{\vec{k}_i^2 + m_i^2}$

# Point Form Spectator Model

- Point form spectator model

$$\begin{aligned} & \langle p'_1, p'_2, p'_3; \sigma'_1, \sigma'_2, \sigma'_3 | \hat{O}_{rd} | p_1, p_2, p_3; \sigma_1, \sigma_2, \sigma_3 \rangle \\ &= 3N \langle p'_1, \sigma'_1 | \hat{O}_{spec} | p_1, \sigma_1 \rangle 2p_{20} \delta(\vec{p}_2 - \vec{p}'_2) 2p_{30} \delta(\vec{p}_3 - \vec{p}'_3) \delta_{\sigma_2 \sigma'_2} \delta_{\sigma_3 \sigma'_3} \end{aligned}$$

- Axial current

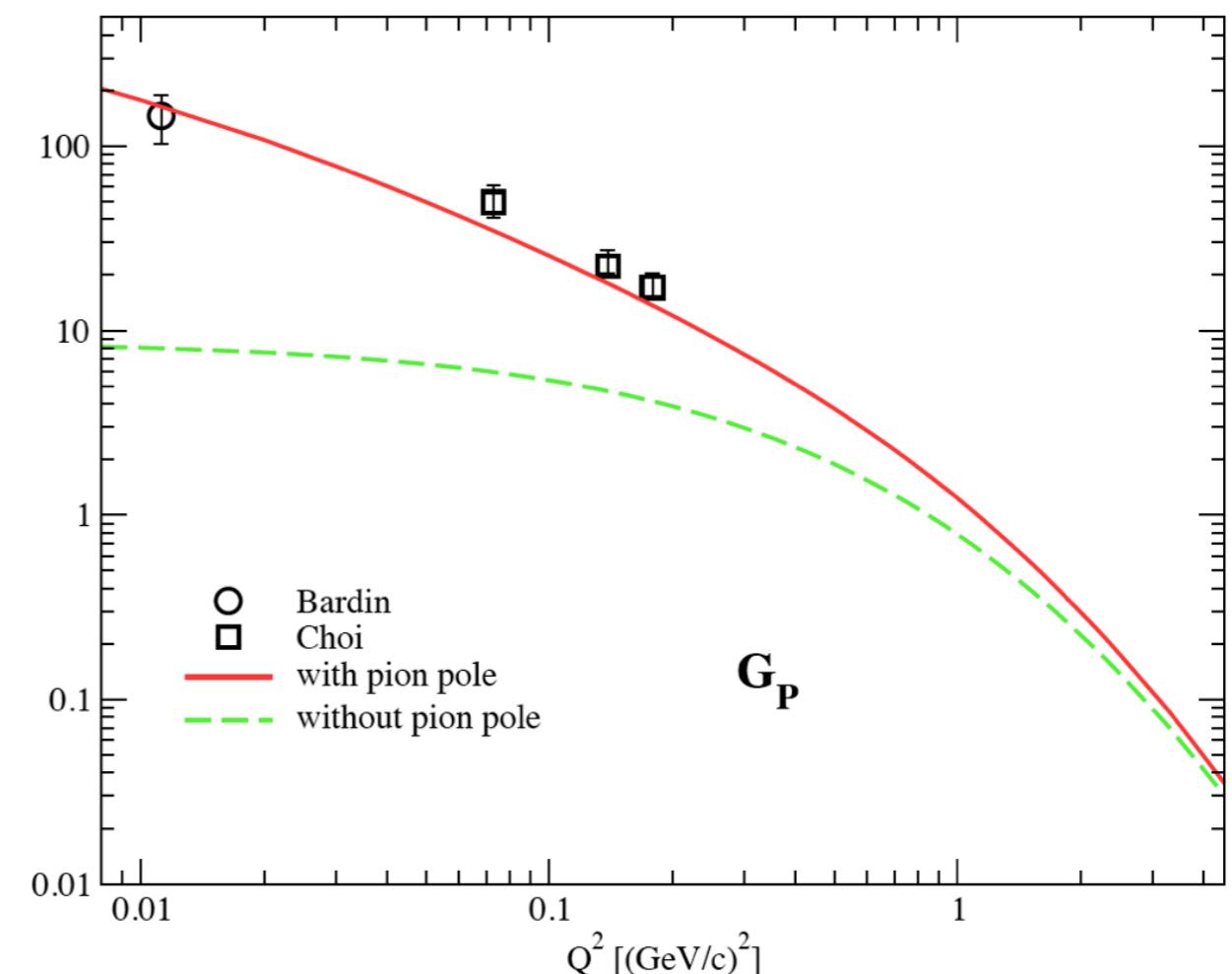
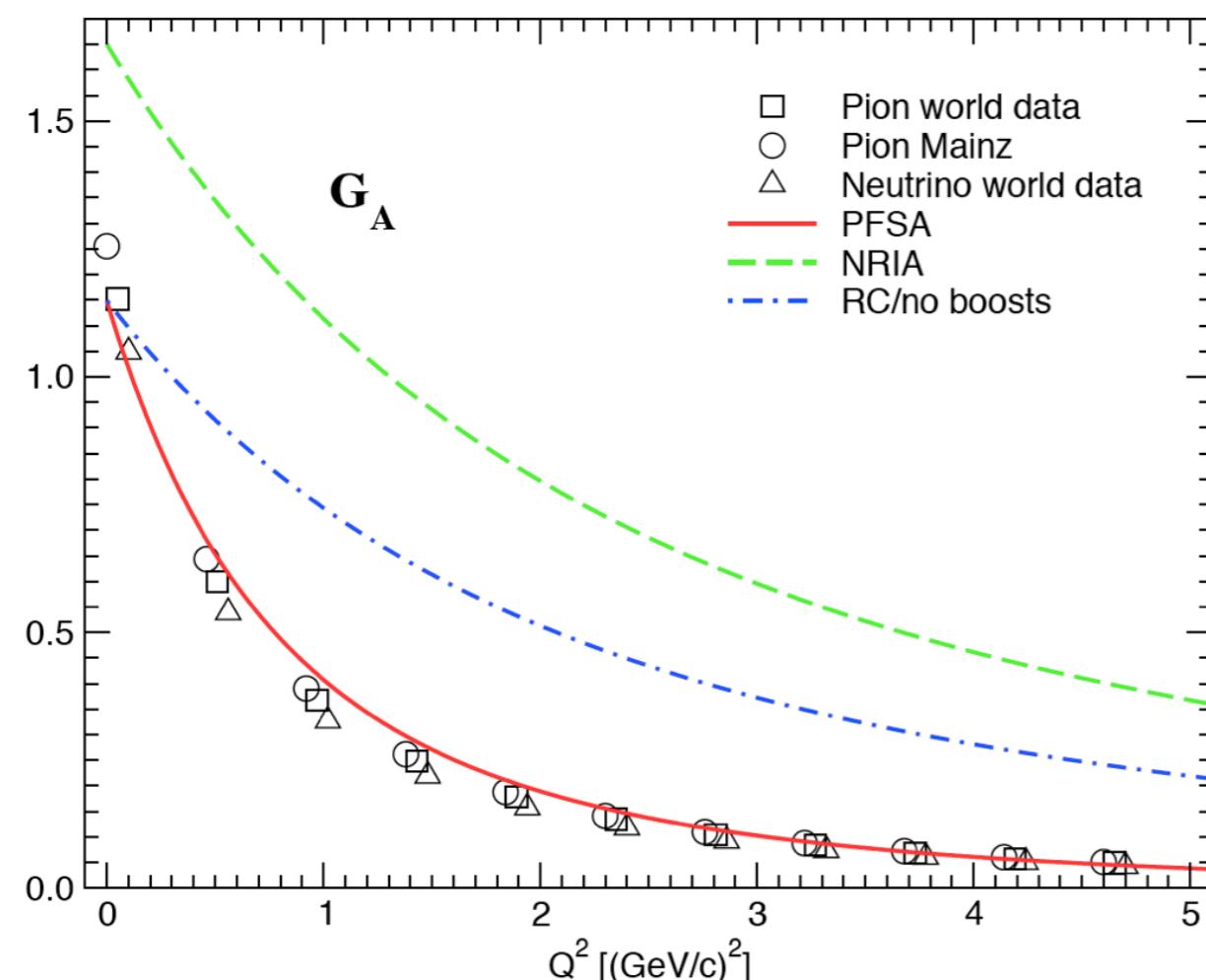
$$\langle p'_1, \sigma'_1 | \hat{A}_{a, spec}^\mu | p_1, \sigma_1 \rangle = \bar{u}(p'_1, \sigma'_1) \left[ g_A^q \gamma^\mu + \frac{2f_\pi}{\tilde{Q}^2 + m_\pi^2} g_{qq\pi} \tilde{q}^\mu \right] \gamma_5 \frac{\tau_a}{2} u(p_1, \sigma_1)$$

$$g_A^q = 1 \quad (\text{pointlike constituent quark})$$

# Results

- Axial form factors of  $N$
- Axial charges of  $N$  and some  $N^*$  relative to other results
- Axial charges of all  $N^*$  resonances below 2 GeV
- Axial form factors of  $N^*$

# Axial and Pseudo-scalar Form Factors of nucleon



$$g_A^{EXP} = 1.2695 \pm 0.0029$$

$$g_A^{psGBE} = 1.15$$

# Comparison

State	$J^P$	EGBE	Lattice QCD <sup>1</sup>	Lattice QCD <sup>3</sup>	GN <sup>2</sup>	NR
$N(939)$	$\frac{1}{2}^+$	1.15	1.10~1.40		1.66	1.65
$N^*(1440)$	$\frac{1}{2}^+$	1.16	—		1.66	1.61
$N^*(1535)$	$\frac{1}{2}^-$	0.02	~0.00		-0.11	-0.20
$N^*(1710)$	$\frac{1}{2}^+$	0.35	—		0.33	0.42
$N^*(1650)$	$\frac{1}{2}^-$	0.51	~0.55		0.55	0.64

1. K. S. Choi, W. Plessas, R. F. Wagenbrunn : Phys. Rev. C 81, 028201 (2010)

2. L. Ya. Glozman, A.V. Nefediev: Eur. Phys. J.A 12, 91 (2001)

3. Toru T.Takahashi and Teiji Kunihiro : Phys. Rev. D 78, 011503 (2008)  
and further lattice QCD calculation

# Axial Charges of Chiral Partners

State	$J^P$	EGBE		psGBE		OGE	
		Mass	$g_A$	Mass	$g_A$	Mass	$g_A$
$N(939)$	$\frac{1}{2}^+$	939	1.15	939	1.15	939	1.11
$N^*(1520)$	$\frac{3}{2}^-$	1524	-0.64	1519	-0.21	1520	-0.15
$N^*(1440)$	$\frac{1}{2}^+$	1464	1.16	1459	1.13	1578	1.10
$N^*(1535)$	$\frac{1}{2}^-$	1498	0.02	1519	0.09	1520	0.13
$N^*(1680)$	$\frac{5}{2}^+$	1689	0.89	1728	0.83	1858	0.70
$N^*(1675)$	$\frac{5}{2}^-$	1676	0.84	1647	0.83	1690	0.80
$N^*(1710)$	$\frac{1}{2}^+$	1757	0.35	1776	0.37	1860	0.32
$N^*(1650)$	$\frac{1}{2}^-$	1581	0.51	1647	0.46	1690	0.44
$N^*(1720)$	$\frac{3}{2}^+$	1746	0.35	1728	0.34	1858	0.25
$N^*(1700)$	$\frac{3}{2}^-$	1608	-0.10	1647	-0.50	1690	-0.47

# Classification

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 $(LS)J^P$ 


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$(0\frac{1}{2})\frac{1}{2}^+$	$N(939)^{100}$	$\Lambda(1116)^{100}$	$\Sigma(1193)^{100}$	$\Xi(1318)^{100}$
$(0\frac{1}{2})\frac{1}{2}^+$	$N(1440)^{100}$	$\Lambda(1600)^{96}$	$\Sigma(1660)^{100}$	<b><math>\Xi(1690)^{100}</math></b>
$(0\frac{1}{2})\frac{1}{2}^+$	$N(1710)^{100}$		$\Sigma(1880)^{99}$	
$(1\frac{1}{2})\frac{1}{2}^-$	$N(1535)^{100}$	$\Lambda(1670)^{72}$	<b><math>\Sigma(1560)^{94}</math></b>	
$(1\frac{3}{2})\frac{1}{2}^-$	$N(1650)^{100}$	$\Lambda(1800)^{100}$	<b><math>\Sigma(1620)^{100}</math></b>	
$(1\frac{1}{2})\frac{3}{2}^-$	$N(1520)^{100}$	$\Lambda(1690)^{72}$	$\Sigma(1670)^{94}$	$\Xi(1820)^{97}$
$(1\frac{3}{2})\frac{3}{2}^-$	$N(1700)^{100}$		<b><math>\Sigma(1940)^{100}</math></b>	
$(1\frac{3}{2})\frac{5}{2}^-$	$N(1675)^{100}$	$\Lambda(1830)^{100}$	$\Sigma(1775)^{100}$	<b><math>\Sigma(1950)^{100}</math></b>

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 $(LS)J^P$ 


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$(0\frac{3}{2})\frac{3}{2}^+$	$\Delta(1232)^{100}$	$\Sigma(1385)^{100}$	$\Xi(1530)^{100}$	$\Omega(1672)^{100}$
$(0\frac{3}{2})\frac{3}{2}^+$	$\Delta(1600)^{100}$	<b><math>\Sigma(1690)^{99}</math></b>		
$(1\frac{1}{2})\frac{1}{2}^-$	$\Delta(1620)^{100}$	<b><math>\Sigma(1750)^{94}</math></b>		
$(1\frac{1}{2})\frac{3}{2}^-$	$\Delta(1700)^{100}$			

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K. S. Choi, W. Plessas, R. F. Wagenbrunn : Phys. Rev. D 82, 014007 (2010)

T. Melde, W. Plessas and B. Sengl: Phys. Rev. D 77, 114002 (2008) and Particle Data Group (PDG) 2010

# Classification

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$(0\frac{1}{2})\frac{1}{2}^+$	$N(1440)^{100}$	$\Lambda(1600)^{96}$	$\Sigma(1660)^{100}$	$\Xi(\mathbf{1690})^{100}$
$(0\frac{1}{2})\frac{1}{2}^+$	$N(1710)^{100}$	<b>0.02</b>	$\Sigma(1880)^{99}$	
$(1\frac{1}{2})\frac{1}{2}^-$	$N(1535)^{100}$	$\Lambda(1670)^{72}$	$\Sigma(\mathbf{1560})^{94}$	-0.15
$(1\frac{3}{2})\frac{1}{2}^-$	$N(1650)^{100}$	$\Lambda(1800)^{100}$	$\Sigma(\mathbf{1620})^{100}$	
$(1\frac{1}{2})\frac{3}{2}^-$	$N(1520)^{100}$	$\Lambda(1690)^{72}$	$\Sigma(1670)^{94}$	$\Xi(1820)^{97}$
$(1\frac{3}{2})\frac{3}{2}^-$	$N(1700)^{100}$			$\Sigma(\mathbf{1940})^{100}$
$(1\frac{3}{2})\frac{5}{2}^-$	$N(1675)^{100}$	$\Lambda(1830)^{100}$	$\Sigma(1775)^{100}$	$\Sigma(\mathbf{1950})^{100}$

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 $(LS)J^P$ 


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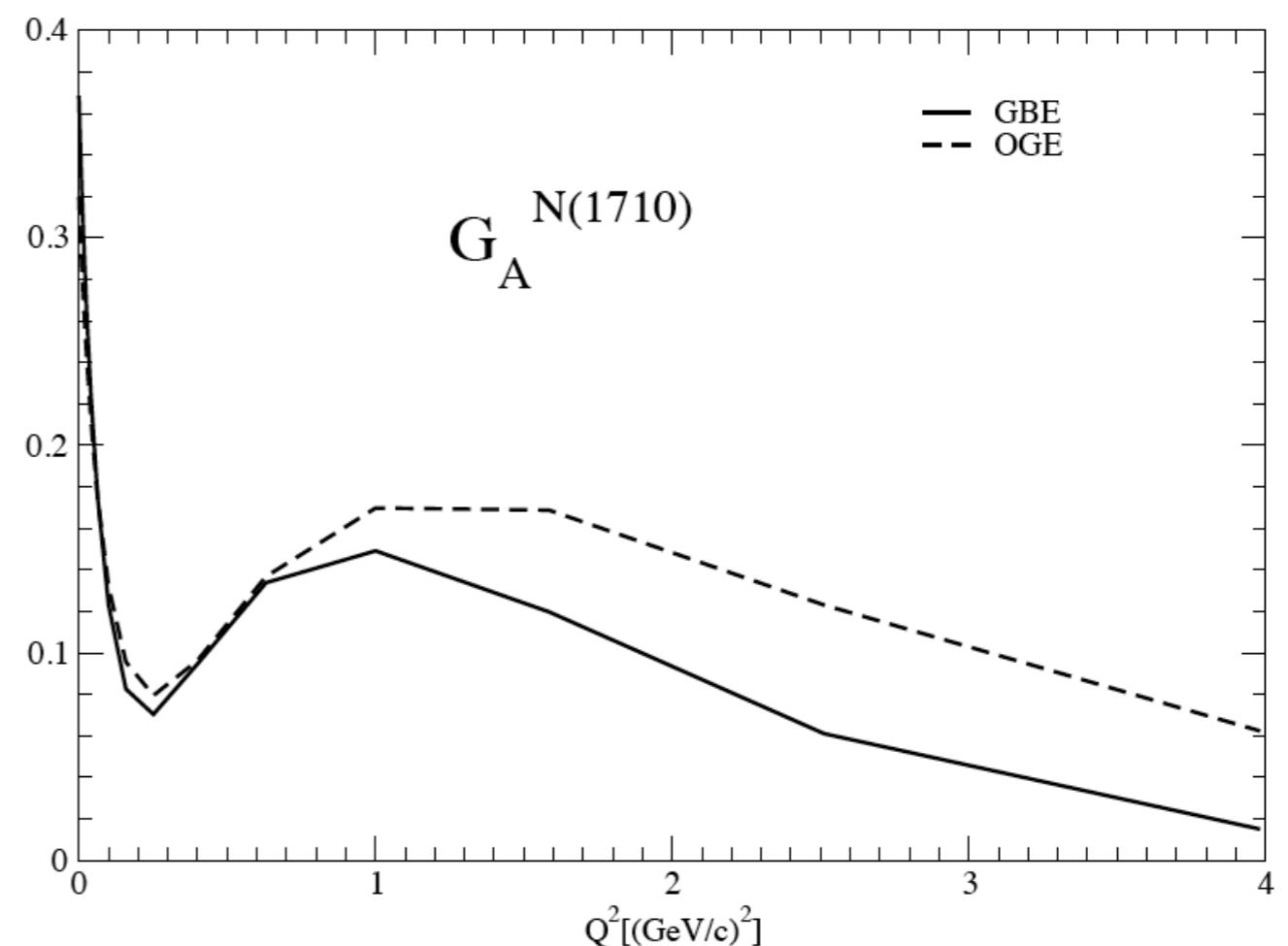
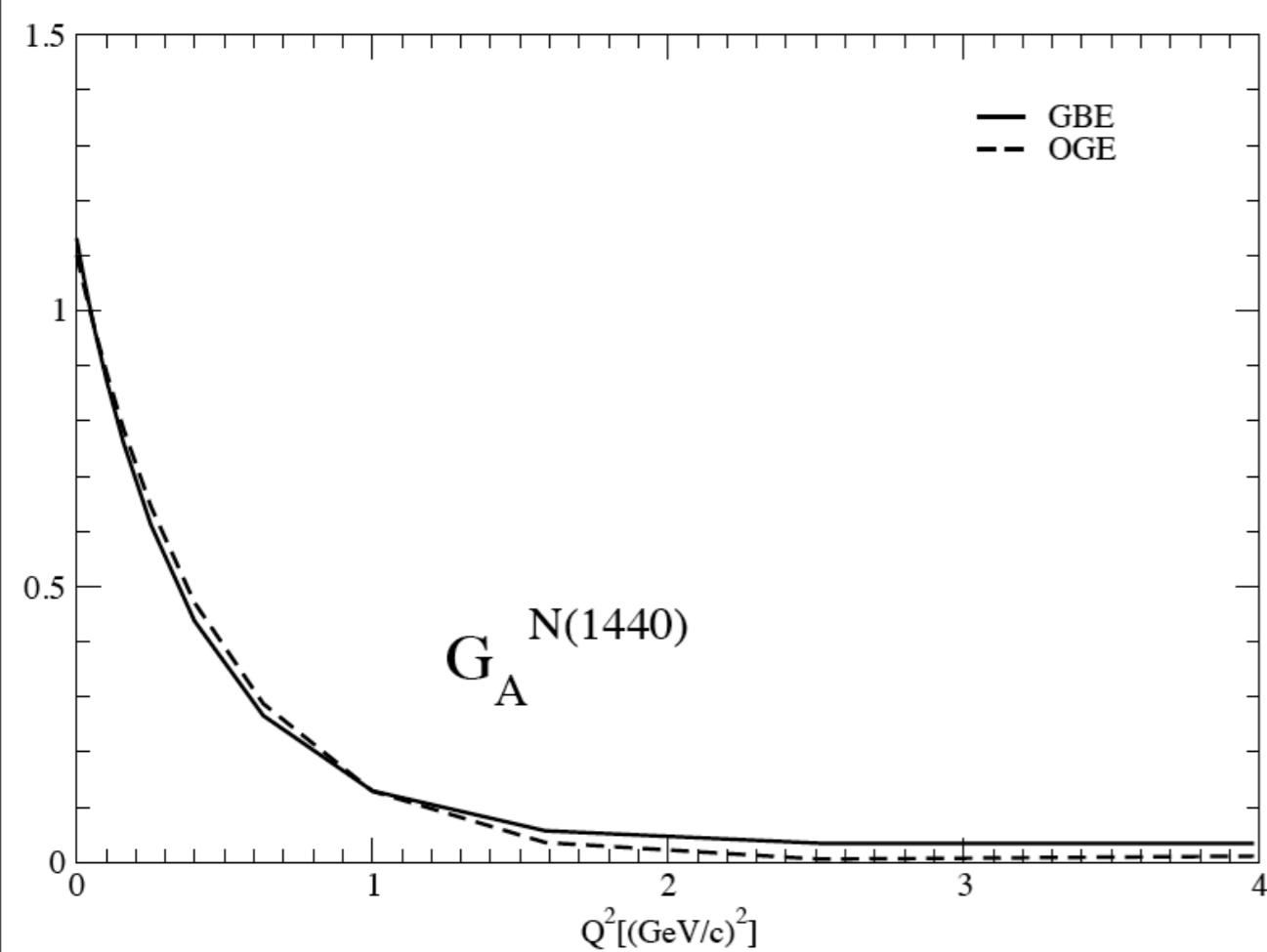
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$(0\frac{3}{2})\frac{3}{2}^+$	$\Lambda(1600)^{100}$	$\Sigma(\mathbf{1690})^{99}$		
$(1\frac{1}{2})\frac{1}{2}^-$	$\Delta(1620)^{100}$	$\Sigma(\mathbf{1750})^{94}$	-0.08	
$(1\frac{1}{2})\frac{3}{2}^-$	$\Delta(1700)^{100}$		<b>-0.76</b>	

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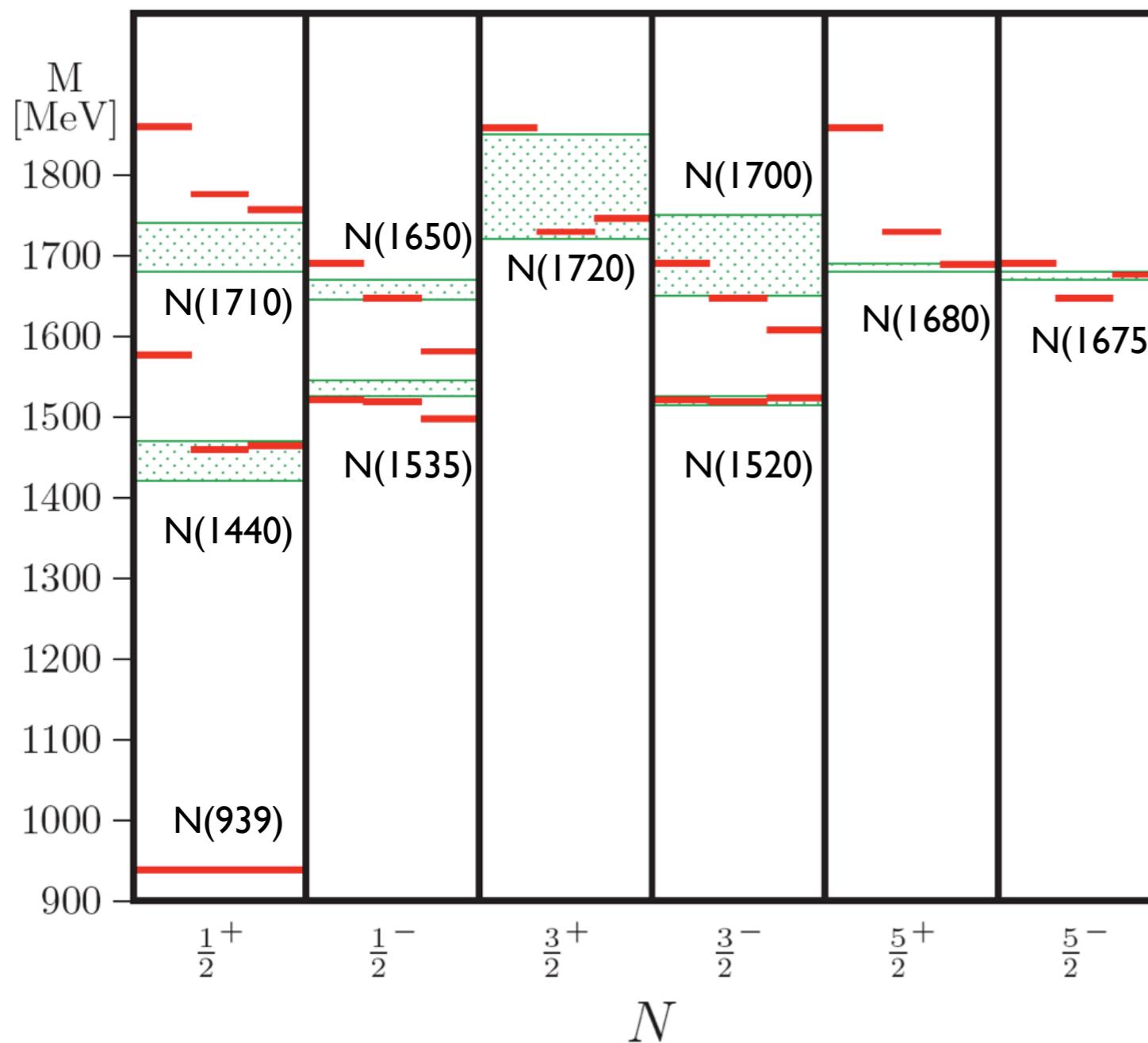
# Axial Form Factors of $N(1440)$ and $N(1710)$



## Summary

- Performed a consistent study of  $g_A$  for N and N\* within the RCQMs.
- Covariant predictions of RCQMs show generally considerable relativistic effects and describe experimental data well.
- In cases, where no experiments exist, there is reasonable agreement with lattice QCD results.
- The issue regarding chiral symmetry restoration remains open; the sizes of the  $g_A$  for baryon resonances are consistent with a recent classification into flavor multiplets.
- It will be interesting to extend these relativistic studies to electromagnetic and weak transition form factors.

# Nucleon Spectra



Left levels - OGE : L.Theussl, R.F.Wagenbrunn, B. Desplanques, and W.Plessas: Eur. Phys. J. A 12, 91 (2001)

Middle levels - psGBE : L.Ya. Glozman, W. Plessas, K. Varga, and R.F.Wagenbrunn: Phys. Rev. D 58, 094030 (1998)

Right levels - EGBE : K.Gantschnig, R.Kainhofer, W. Plessas, B. Sengl, and R.F.Wagenbrunn: Eur. Phys. J. A 23, 507 (2005)

Green zones - Experimental data in PDG (2010)